Interoffice

National Radio Astronomy Observatory

Charlottesville, Virginia

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To: M. Davis and Engineering

From: S. von Hoerner

Subject: Non-Homologous Deformations of the 300-ft.

In his memo of April 2, 1969, M. Davis has plotted the aperture efficiency n of the 300-ft telescope as a function of zenith angle ϕ , as measured at $\lambda = 21.4$ cm. His Fig. 3 shows three curves n (ϕ): (a) <u>1962-66</u>, before readjustment; (b) <u>1967-69</u>, after readjustment; showing the single measured points for both curves; and (c) a <u>prediction</u> for the new surface. By comparison of two wavelengths (21.4 and 40 cm) he also finds the efficiency n₀ for $\lambda \rightarrow \infty$ as (a) n₀ = .67 and (b) n₀ = .59 and he adopts (**6**) n₀ = .63.

In the following, I make a best-fit of a theoretical formula to (a) and (b) and derive a somewhat different prediction (c), using the same three values n_0 as M. Davis. Let a telescope be adjusted without gravity to a perfect paraboloid. Then, with gravity, call ΔH_1 the rms deviation of the surface from a bestfit paraboloid in zenith position, and ΔH_2 likewise in horizon position. These two parameters fully describe the gravitational effects.

If a telescope is adjusted to a perfect paraboloid at zenith angle θ , and then observes at zenith angle ϕ , the deviation ΔH from a best-fit paraboloid is

$$\Delta H = \sqrt{\Delta H_1^2 (\cos \phi - \cos \theta)^2 + \Delta H_2^2 (\sin \phi - \sin \theta)^2}$$
(1)

If the surface itself has an rms error σ_0 , the total rms deviation from a paraboloid then is

$$\sigma = \sqrt{\Delta H^2 + \sigma_0^2}$$
 (2)

and the aperture efficiency is

$$\eta = \eta_0 e^{-(4\pi\sigma/\lambda)^2}$$
(3)

Regarding ΔH_1 and ΔH_2 , it would be interesting to see whether or not there is a difference before and after the readjustment, since this was connected with some strengthening of the back-up structure (mainly the wheel); but the data cover too small a range in ϕ and scatter too much for this purpose. There seems to be a small improvement (about 20%), but in the following we neglect the difference and adopt the same ΔH_1 and ΔH_2 for (a), (b) and (c).

Since the second term in (1) is always much larger than the first one for the range of ϕ covered, we have a large uncertainty for ΔH_1 ; but this does not

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effect prediction (c) very much, for the same reason. The best-fitting values and their estimated mean errors are

$$\Delta H_1 = 18.5 \pm 5.0 \text{ mm}$$
 (4)

$$\Delta H_2 = 7.5 \pm 1.5 \, \text{mm} \tag{5}$$

and their ratio is

$$g = \Delta H_2 / \Delta H_1 = 0.405 \tag{6}$$

In a reproduction of M. Davis' Fig. 3, I have entered the points calculated with equation (1), using parameters (4) and (5) and adopting the same values n_0 as M. Davis. The agreement with the measured points is certainly within the scatter of the data.

Next, prediction (c) is plotted for a new surface adopting σ_0 = 4 mm as M. Davis did. This new prediction gives smaller efficiencies for large ϕ than the old one.

Finally, if the best adjustment angle θ is defined by the demand

$$\Delta H_{-30^{\circ}} = \Delta H_{+60^{\circ}} \tag{7}$$

see Fig. 2, then one obtains from (1) and (6) that $\theta = 29.3^{\circ}$ or roughly

$$\theta = 30^{\circ} \tag{8}$$

With the available data, the uncertainty of g is rather large. With a probable error range of

$$0.29 < g < 0.58$$
 (9)

we find from Fig. 2

$$22^{\circ} < \theta < 36^{\circ}$$
 (10)

Prediction (c) also has a large uncertainty. From (4) and (5) we obtain, for example, the probable error ranges

$$10.08 < \sigma (60^{\circ}) < 14.38$$
 (11)

and

$$0.324 \le n (60^\circ) \le 0.465.$$
 (12)

SvH:jab



