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A preliminary value for  $B_1$  has been given by Wade (1). This report gives a new determination of this parameter using all the available data for baseline 2. In addition, the value of  $B_3$  during the observing period November 1964 - January 1965 is determined.

#### I. Procedure.

The procedure has been derived by Wade (1). The phase output of the calibration program is given by

$$K = \Delta B_3 - B_1 \Delta \delta \cos \delta - \Delta B_1 \sin \delta$$
 (1)

where  $\Delta\delta$  is the error in the assumed declination  $\delta$  of the source

 $\Delta B_1$ ,  $\Delta B_3$  are the errors in the assumed values of  $B_1$  and  $B_3$ .

The difference in mean phase of two sources of known declinations  $\delta_{1}$  and  $\delta_{2}$  is therefore

$$\Delta \mathbf{K} = \mathbf{K}_1 - \mathbf{K}_2 = \Delta \mathbf{B}_1 \ (\sin \delta_2 - \sin \delta_1)$$

The phase outputs of the calibration program are ambiguous by multiples of  $2\pi$ ; to allow for this, the expression used in the solution of  $\Delta B_{\perp}$  includes an integer N = 0,  $\pm$  1,  $\pm$  2, etc. Thus

$$\Delta B_{1} = \frac{\Delta K}{\sin \delta_{2} - \sin \delta_{1}} + \frac{N}{\sin \delta_{2} - \sin \delta_{1}}$$
(2)

Finally, for a number of sources an incorrect declination had been assumed in the fringe reduction program. For these, the quantity  $B_1 \Delta \delta \cos \delta$  was included in the K term, as indicated in equation (1).

# II. The Mean Value of B<sub>1</sub> During the Observing Period.

The best data available for the determination of  $B_1$  are presented in Table I. In general, only those observations which covered at least six hours of hour angles, and whose solutions have errors of less than 10° in phase were used. An exception was made for the solution 3C 119 - 3C 286, where two of the three tapes are of lower quality; the value of  $B_1$  thus found was given lower weight in the final mean. Column 6 of Table I gives the value of  $\Delta B_1$  computed from equation (2) with N = 0. Various values of the increment  $\frac{N}{\sin \delta_2 - \sin \delta_1}$  were then

applied to each source until the value of  $B_1$ ', which is consistent with all pairs of sources, was found. This solution is given in column 7 of Table I. The pair 3C 147 - 3C 380 was used in the determination of the approximate value of  $B_1$ ', since an uncertainty of  $2\pi$  corresponds to an increment of 78 wavelengths. The widely separated sources, e.g. 3C 48 - 3C 147 then give the accurate value of  $B_1$ '.

	The Determi	nation of B <sub>l</sub>	Baseline 2 N	ov. 1964 -	- Jan. 190	55
		$K_1 - K_2$	1			Best Solution
Sources	Tape No.	(circles)	sin $\delta_1$ -sin $\delta_2$	Bl	ΔΒι	B <sub>l</sub> '
(1) 3C 147	49	+0.5186	78.446	5030.00	-40.68	5070.7
(2) 3C 380	* 50	+0.6942	78.418	5030.00	-54.44	5084.4
	65	+0.5386	78.285	5030.00	-42.16	5072.2
	66	+0.0408	78.269	5077.72	- 3.19	5080.9
	67	-0.1425	78.261	5077.72	+11.15	5066.6
	*Phase o	utput correc	ted for an erro	$r \Delta \delta = -20$	)"6 (-119	°)
	Mean B <sub>l</sub>	' = 5074.6 <u>+</u>	3.3 (m.e.)	1	1	I
(1) 3C 48	34	+0.4383	29.437	5030.00	-12.90	5072.34
(2) 3C 286	60	. 4353	29.394	5030.00	-12.80	5072.19
	61	.4714	29.391	5030.00	-13.86	5073.25
	70	+0.8628	29.383	5077.72	-25.35	5073.69
	Mean B <sub>l</sub>	' = 5072.87	<u>+</u> 0.36 (m.e.)			
(1) 3C 147	20	+0.5442	4.549	5030.00	- 2.48	5073.42
(2) 3C 48	21	.4703	4.549	5030.00	- 2.14	5073.08
	22	.4158	4.549	5030.00	- 1.89	5072.83
	46	. 5589	4.549	5030.00	- 2.54	5073.48
	Mean B <sub>l</sub>	' = 5073.20	<u>+</u> 0.15 (m.e.)			
(1) 3C 216	* 32	-0.0861	7.241	5030.00	+ 0.62	5072.82
(2) 3C 48	45	-0.0742	7.242	5030.00	+ 0.54	5072.92
	46	+0.0053	7.242	5030.00	- 0.04	5073.49
	Mean B <sub>l</sub>	' = 5073.08	+ 0.21 (m.e.)	•	•	1
	*Phase of	utput correc	ted for an erro	r Δδ -2" ( 1	(-14°)	1
(1) 3C 119	* 29	0.6072	6.534	5030.00	- 3.97	5073.17
(2) 3C 286	33	0.4978	6.534	5030.00	- 3.25	5072.46
	34**	0.5567	6.534	5030.00	- 3.64	5072.84
	Mean B <sub>l</sub>	' = 5072.83	+ 0.18		•	•
	*Phase co	prrected for	Δδ – 8"7 (–57°	) **Given	double we	eight in mean
(1) 3C 38C	* 22	-0.0053	4.828	5030.00	+ 0.03	5073.43
(2) 3C 48	32-33	-0.0222	4.829	5030.00	+ 0.11	5073.35
	33-34	+0.0497	4.829	5030.00	- 0.24	5073.50
	34-35	-0.0694	4.829	5030.00	+ 0.34	5073.13
	35-36	+0.0644	4.829	5030.00	- 0.31	1 5073.77
	Mean B <sub>l</sub>	' = 5073.48	+ 0.11			
	*Phase c	orrected for	Δδ -20"6 (-119	<b>~</b> )		

Overall mean (weights 0, 1, 2, 2, 1, 2) 5073.15 + 0.09

TABLE I

The mean value of  $B_1$  for the period is

$$B_1 = 5073.15 + 0.09 (m.e.)$$

Here the individual determinations have been given weights 0, 1, 2, 2, 1 and 2 respectively. Since the mean value of  $B_2$  was

 $B_{2} = 12493.544 \pm 0.005,$   $\frac{D}{\lambda} = 1.348427 \times 10^{+4} \text{ (Since } B_{1}^{2} + B_{2}^{2} = 1.818254 \times 10^{8}\text{)}$   $\cos d = 0.926527_{3}$   $\sin d = 0.376227_{3}$   $d = 22^{\circ} 06^{\circ} 01^{\circ\circ}$ 

It has been shown previously that  $D/\lambda$  varied during the observing period. It is not possible to see the effect on  $B_1$  itself, because of the large uncertainties in the individual determinations. However, an approximate correction, found by scaling the error  $\Delta B_2$  by the ratio  $\frac{B_1}{B_2}$ = 0.40606 can be applied to the mean value of  $B_2$ . The adopted values

= 0.40606, can be applied to the mean value of  $B_1$ . The adopted values of the error  $\Delta B_1$  as a function of tape number are given in Table II.

The Error	$\Delta B_1$ as a Function	of Tape Number
Tape No.	Error in B <sub>l</sub> Due to LO Drift Along	Total Error ∆B <sub>l</sub>
20 21 22 23 24 25 26	$ \begin{array}{r} -0.02 \\ -0.02 \\ -0.03 \\ -0.02 $	-43.17 -43.17 -43.18 -43.17 -43.17 -43.17 -43.17
27 28 29	-0.02 -0.02 -0.02	+ 4.55 -43.17 -43.17
31 32 33	-0.01 -0.01 -0.01 -0.01	-43.16 -43.16 -43.16 -43.16
34 35 36 37 38	-0.01 -0.01 -0.01 -0.01	-43.16 -43.16 -43.16 -43.16 -43.16
39	0.00	-43.13

TABLE II

TABLE II (Continued)					
Tape No.	Error in B <sub>l</sub> Due to LO Drift Along	Total Error ∆B <sub>l</sub>			
40	0.00	-43.15			
41	0.00	-43.15			
42	0.00	-43.15			
43	0.00	-43.15			
44	0.00	-43.15			
45	0.00	-43.15			
46	+0.01	-43.14			
47	+0.01	-43.14			
48	+0.01	-43.14			
49	+0.01	-43.14			
50	+0.01	-43.14			
51	+0.01	-43.14			
52	+0.01	-43.14			
56	+0.01	-43.14			
57	+0.01	-43.14			
58	+0.01	-43.14			
59	+0.01	-43.14			
60	+0.01	-43.14			
61	+0.01	-43.14			
62	+0.01	+ 4.58			
63	+0.02	+ 4.59			
64	+0.02	+ 4.59			
65	+0.02	-43.13			
66	+0.02	+ 4.59			
67	+0.02	+ 4.59			
68	+0.02	+ 4.59			
69	+0.02	+ 4.59			
70	+0.02	+ 4.59			
71	+0.02	+ 4.59			
72	+0.02	+ 4.59			

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# III. The Value of $\Delta B_3$ During the Observing Period.

The parameter  $\Delta B_3$  (equal to -  $B_3$  for all tapes) is obtained from equation (1) using the values of  $\Delta B_1$  from Table II. Only the observations of the calibrators which covered six hours of hour angle and for which the rms scatter in phase is less than 10° have been included. The values of  $\Delta B_3$  are given in Table III, along with the uncertainties in K. The uncertainty in  $B_1$  of 0.09 wavelengths produces an uncertainty in  $\Delta B_3$  as given at the bottom of Table III. It should be noted, however, that these uncertainties are coupled for the various sources, and are therefore essentially an uncertainty in the zero point of the plot. Figure 1 gives the value of  $\Delta B_3$  as a function of tape number for all of the calibration sources. This figure is probably the best summary of the long-term phase stability of the system during the observing period November 1964 - January 1965.

Tape	<b>3</b> C 48	3C 147	3C 196	3C 286	3C 380
20 21 22 23 24 25	$\begin{array}{r} + 51 + 3 \\ - 01 + 2 \\ - 05 + 1 \end{array}$	$\begin{array}{r} + 72 + 1 \\ - 07 + 3 \\ - 31 + 3 \\ + 09 + 3 \\ - 94 + 1 \\ - 228 + 1 \end{array}$			+ 16 <u>+</u> 2
26 27 28					-222 <u>+</u> 3
29 30 31	-260 <u>+</u> 6	-143 + 1	-151 <u>+</u> 3	-197 <u>+</u> 10	-161 <u>+</u> 2
32 33 34 35 36 37	$\begin{array}{r} -160 + 4 \\ -173 + 2 \\ -162 + 2 \\ -184 + 3 \end{array}$	-135 + 2	-181 + 5 -157 + 4 -129 + 7	-156 + 7 -162 + 6	$ \begin{array}{r} -144 + 9 \\ -143 + 5 \\ -159 + 4 \\ -138 + 3 \\ -157 + 4 \end{array} $
38 39 40	-222 + 2		-174 <u>+</u> 11		-113 + 3 -158 + 3
41 42 43 44	-208 <u>+</u> 7 - 04 <u>+</u> 4	- 46 <u>+</u> 1	- 82 <u>+</u> 3		- 57 <u>+</u> 2
45 46 47 48 49 50 51	-39 + 1 + 45 + 1	$ \begin{array}{r} -26 + 3 \\ +73 + 1 \\ +80 + 1 \\ +63 + 2 \\ +71 + 2 \\ +98 + 4 \\ +97 + 4 \end{array} $			+ 82 + 2 + 47 + 3 + 57 + 3
52 56 57 58 59 60	$\begin{array}{r} + 52 + 1 \\ + 54 + 10 \\ + 47 + 4 \\ + 24 + 2 \end{array}$			$+108 \pm 2$ + 76 \pm 3 + 35 \pm 1	

TABLE III Summary of Values of  $\Delta B_3$ 

			(001102110207		
Tape	3C 48	3C 147	3C 196	3C 286	3C 380
61 62 63 64 65 66 67 68 69	$\begin{array}{r} + 44 + 1 \\ + 25 + 2 \\ + 35 + 2 \end{array}$	$\begin{array}{r} + & 09 + 2 \\ + & 10 + 1 \\ + & 27 + 1 \\ + & 09 + 2 \\ + & 15 + 5 \end{array}$		+ 28 + 3 + 30 $+ 3$	$\begin{array}{r} + 22 + 4 \\ + 14 + 4 \\ + 40 + 3 \\ + 20 + 3 \\ \end{array}$
70 71 72 Uncertain to error	+ 21 + 2 + 26 + 2 ty due in ΔB <sub>1</sub> 18°	+ 22 <u>+</u> 1 25°	24°	+ 14 <u>+</u> 1 17°	24°

TABLE III (Continued)

### IV. The determination of Source Declination from the Mean Phase.

The phase drift method for the determination of source declination becomes inaccurate for sources near the equator, since the correction determined is a function of the cosecant and cotangent of the declination. The declination correction can, in principle, be determined much more accurately from the observed phase, since, from equation (1)

$$\Delta \delta = -\frac{K - \Delta B_3 + \Delta B_1 \sin \delta}{B_1 \cos \delta}$$
(3)

Thus, for example, an uncertainty in the quantity  $K - \Delta B_3$  of 20° leads to an uncertainty of only 2" in the declination correction for a source within 5° of the equator.

There were three such sources included in the position program run during this observing period. The corrections in declination, as determined using equation (3), are given in Table IV. Table V gives the comparison between the declinations obtained by this method and those of the phase drift program. The errors are smaller, as expected, and moreover, the position of 3C 298 is in much better agreement with the optical position. The position of 3C 459 is very uncertain, since the one observation was poor.

Source Declination from Mean Phase							
Source	Tape	K -∆B <sub>3</sub> (circles)	ΔB <sub>l</sub> sin δ	$\frac{1}{B_1 \cos \delta}$	Primary Value Δδ radians	"Best" Value Δδ (secs)	Mean ∆ô
273 298 459	56 57 58 27 33 34 40	-0.128 -0.010 -0.273 +0.486 +0.091 +0.027 +0.331	-1.692 -1.692 -1.692 +0.526 -4.991 -4.991 -2.934	$1.973 \times 10^{-4}$ 1.973 × 10 <sup>-4</sup> 1.973 × 10 <sup>-4</sup> 1.985 × 10 <sup>-4</sup> 1.985 1.985 1.985 1.976 × 10 <sup>-4</sup>	+1.62 x $10^{-4}$ 1.38 1.90 -0.024x $10^{-4}$ -1.77 -1.91 +1.19	$\begin{array}{r} - 7.3 \\ -12.1 \\ - 1.4 \\ - 0.5 \\ - 4.5 \\ - 1.5 \\ -16.2 \end{array}$	-6.9 <u>+</u> 3.1 -2.2 <u>+</u> 1.2

TABLE IV unce Declination from Mean Pha

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TABLE V

Source	δ <sub>R</sub> - δ <sub>o</sub> Phase Drift	δ <sub>R</sub> - δ <sub>o</sub> Mean Phase
3C 273	+ 6.6 <u>+</u> 7.4	+ 6.9 <u>+</u> 3.1
298	-12.1 <u>+</u> 4.4	+ 2. <b>2</b> <u>+</u> 1.2
459	+9?	+15.5(+25.3)

It is clear that the phase characteristics of the interferometer can be measured accurately enough to allow meaningful positions to be determined from the observed phase of the source. This is most important, since it means that sources near the equator can be measured with the same accuracy as is achieved elsewhere with the phase-drift method.

## References

 Wade, C. M. "The Value of B<sub>1</sub> for Interferometer Baseline 2", NRAP Report January 1965.





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