## NATIONAL RADIO ASTRONOMY OBSERVATORY



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## MEMO To: John Webber

From:	Skip Thacker		
Date:	January 14, 1999		

### Subject: MMA Total Power: Correlation Receivers

The MMA total power mode cries out for a correlation Receiver of the type shown in figure 7-21 of Krause (see attached references). An example block diagram of how a 100 GHz system could be constructed using HEMT amplifiers is attached as figure 1. Note the similarity to block diagram of Predmore figure 1 (see attached references). Paralleling Predmore's analysis except substituting (G1 + deltaG1) for his G1 and G2 + deltaG2 for his G2 where my G1 is the average value of the gain and deltaG1 is a zero mean RV that is uncorrelated with any of the T's and G's. This is to show that the correlation receiver is insensitive to gain variation.

The really neat comparison is to look at figure 5.3.4(a) in the MMA project book (see attached) and compare that to Predmore figure 1. All you have to do to make this into a correlation receiver that is insensitive to amplifier gain fluctuations is to put the IF amplifier in front of the 180 degree hybrid. One could possibly eliminate the hybrid and go straight into the correlator and do the combining in software. This of course doubles the number of IF amplifiers that you need, but this is a small price to pay particularly if these are MMIC amplifiers. One also can phase switch this receiver by putting a 180 degree phase shift in one of the LO's of the SIS mixer.



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When uncorrelated noise voltages due to the noise temperatures  $T_{s1}$ and  $T_{s2}$  are multiplied, the product has an average value zero (zero d-c output voltage from the multiplier). As in the detector of the total-power receiver the uncorrelated noise-voltage components from both receivers beat with each other in the multiplier, resulting in a low-frequency, fluctuating noise-voltage output. Referring to (7-11) for the total-power receiver, the fluctuating noise power from the integrator will be

$$W_{\rm LF} = G_{\rm LF} C' k T_{s1} k T_{s2} \,\Delta\nu_{\rm HF} \,\Delta\nu_{\rm LF} \tag{7-38}$$

A discrete source produces correlated signal-noise powers  $k \Delta T \Delta \nu_{\rm HF}$ at the receiver inputs and corresponding IF output voltages with amplitudes



Fig. 7-19. Interferometer with correlation receiver.

proportional to  $\sqrt{k \Delta T \Delta \nu_{\rm HF}}$ . If both are in phase, the multiplier output d-c voltage is proportional to  $k \Delta T \Delta \nu_{\rm HF}$ . If the phase angle between them is  $\phi$ , the d-c output voltage is equal to  $k \Delta T \Delta \nu_{\rm HF} \cos \phi$ . The signal output power from the integrator is, from (7-12),

$$W = G_{\rm LF} C' (k \,\Delta T \,\Delta\nu_{\rm HF})^2 \cos^2 \phi \tag{7-39}$$

The sensitivity of the correlation receiver is obtained by putting  $W = W_{\rm LF}$  from (7-38) and (7-39), or

$$\Delta T_{\min} = \frac{1}{\cos \phi} \sqrt{\frac{T_{s1} T_{s2} \,\Delta \nu_{\rm LF}}{\Delta \nu_{\rm HF}}} \tag{7-40}$$

or if  $T_{s1} = T_{s2} = T_{sys}$  and  $t_{LF} = 1/2\Delta\nu_{LF}$ 

$$\Delta T_{\min} = \frac{1}{\sqrt{2}\cos\phi} \frac{T_{\text{sys}}}{\sqrt{\Delta\nu_{\text{HF}} t_{\text{LF}}}}$$
(7-41)

The sensitivity of a correlation receiver is hence  $2\sqrt{2}$  times better than the sensitivity of a Dicke receiver with the same system noise temperature and one antenna.

Because only correlated noise voltages give a d-c output, voltagegain instabilities will not affect the sensitivity of the correlation receiver. Gain variation will change only the calibration of the receiver. However, random phase variations in the amplifiers of the predetection sections are undesirable (Fujimoto, 1964). For the same reason scintillations in the

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ionosphere will reduce the sensitivity. One of the advantages of the correlation receiver is that there is no switch and, hence, no extra losses between the antenna and the receiver, which means that the noise temperature of the receiver will be lower.



Fig. 7-20. Phase-switched receiver

The correlation principle is also applied in the phase-switched interferometer (Ryle, 1952) in Fig. 7-20. The IF signal of one receiver goes through a phase-reversing switch, which is operated at the frequency  $\nu_M$ . If signals  $v_1$  and  $v_2$  are uncorrelated, the switching will have no effect on the square-law-detector output. When  $v_1$  and  $v_2$  contain correlated components, the detector output is different for  $v_1 + v_2$  and for  $v_1 - v_2$ . This means that the detector output varies at the frequency  $\nu_M$  because of the correlated



Fig. 7-21. Correlation receiver.

signal. Assuming that the desired signal is the only correlated signal, it is clear that the sensitivity of the phase-switching receiver is the same as the sensitivity of the simple Dicke receiver using a similar low-frequency section (O'Donnell, 1963).

The correlation technique can be used with one antenna by dividing the output signal from the antenna between two identical receivers (Fig. 7-21). In this case the antenna noise power  $(T_A/2)$  is also correlated in addition to signal noise. Hence, this modification is useful only when  $T_A$ 

# A Continuous Comparison Radiometer at 97 GHz

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Abstract — A continuous comparison radiometer has been implemented at 97 GHz using quasi-optical techniques for local oscillator (LO) injection and realization of a 90° hybrid. Cryogenically cooled Schottky-diode mixers and FET amplifiers give a double-sideband (DSB) system temperature of 250 K. The system is self-calibrating and optimized under computer control. The root-mean-square (rms) fluctuations due to the receiver are less than 0.025 K with a 1-s integration time.

#### I. INTRODUCTION

A HIGH-SENSITIVITY, continuous comparison radiometer has been implemented by the Five College Radio Astronomy Observatory. The continuous comparison or correlation radiometer circuit is over twenty years old [1] and has been used for radio astronomical observations at 74 cm [2], 11 cm [3], and 6 cm [4]. Its use at millimeter wavelengths has been precluded due to the losses and imperfections in the available waveguide components. The extensive use of quasi-optical techniques for beam guidance, local oscillator (LO) injections, and a mbination of the two beams in an input hybrid has sulted in major performance gains for the system deibed here.

The classical problem in radiometry is to distinguish a weak source in the presence of the much greater noise from the receiver and atmosphere and in the presence of fluctuations in receiver gain and atmospheric emission. For an astronomical object such as a quasar, this is accomplished by actually moving the telescope between the source position and an adjacent position at the same elevation at a 0.05–0.10-Hz rate, deflecting the beam by nutating the subreflector at a 0.5–5-Hz rate [5], Dicke switching [6] against a load, or deflecting the beam near the Cassegrain focal plane with a rotating beam chopper at up to 50 Hz [7].

In contrast, the continuous comparison radiometer has no moving parts for switching, so that the complexities of mechanical switching systems mentioned above, and the significant loss and mismatch of a ferrite switch, are avoided. Rather, two beams are continually subtracted to give the desired difference signal with an equivalent time constant which can be short as the reciprocal of the inter-

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mediate frequency (IF) bandwidth. Since the subtraction is performed continuously, fluctuations with very short time scales that are common to both beams are cancelled perfectly. In addition, the subtraction in the radiometer is insensitive to receiver gain changes.

The present system employs a LO frequency of 97 GHz and is sensitive to input signals in both sidebands with a separation from the LO of 4.4 to 5.0 GHz. The radiometer consists of a calibration system, cryogenic front end, and IF and signal processing units. The entire system is automatically optimized, calibrated, and operated by a computer. The root-mean-square (rms) sensitivity due to the receiver is less than 0.025 K with a 1-s integration time.

#### II. BASIC THEORY

The radiometer block diagram is shown in Fig. 1. The two input beams with temperature  $T_A$  and  $T_B$  as well as the 97.3-GHz LO are split by a quasi-optical 90° hybrid and double-sideband (DSB) mixed to 4.7 GHz. At 4.7 GHz, the voltages in the two IF amplifiers are proportional to A - jB and A + jB, where A and B are the voltages of the inputs to the Dewar. The second frequency conversion is from 4.7 to 2.1 GHz using a 6.8-GHz LO. One of the LO lines has a 6-bit computer-controlled phase shifter with a 5.6° resolution to optimize the system sensitivity. Also under computer control are 0-50-dB p-i-n attenuators in each of the 2.1-GHz IF chains. Precision IF 180° hybrid, matched detectors, and an instrumentation operational amplifier are used to multiply the two IF signals.

The resultant response of the system is

$$V_{\text{DIFF}} = C_{\text{DIFF}} (T_A - T_B) \cos \phi \qquad (1)$$

where  $C_{\text{DIFF}}$  is a calibration constant and  $\phi$  is the phase of the LO phase shifter. Thus, the system is insensitive to gain fluctuations and is responsive only to the input temperature difference. A manual delay line with a 0.2-ns range is used to compensate for phase slopes in the cabling and amplifiers. The entire system is automatically optimized and calibrated with a ModComp MODACS computer.

Faris [1] has done an extensive analysis of correlation radiometers including the effects of nonidentical amplifiers, gain fluctuations, and differential phase and delay. Here, the system response will be analyzed for a single IF frequency to investigate the effect of amplitude imbalance in the input 90° hybrid and imperfect balance in the final IF detectors.

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Fig. 1. Radiometer block diagram. The two inputs at temperatures of  $T_A$  and  $T_B$  are split and combined in a quasi-optical hybrid. After mixing at 97 GHz, they are amplified and phase shifted to optimize the response. A precision 180° hybrid and matched detectors are used to "multiply" the two IF signals.

The amplitude of the incoming signal in the upper sideband (USB) and in the lower sideband (LSB) is proportional to the square root of the input temperature. In the following equations, the unprimed quantities are from the USB and primed quantities from the LSB. The complex conjugate is represented by the symbol (\*). The input temperatures in the signal and reference beams are denoted by  $T_A$  and  $T_B$ , respectively, and the two receiver temperatures are  $T_{R1}$  and  $T_{R2}$ . The signal amplitudes are A and A' and the reference amplitudes B and B', in the two sidebands. The amplitudes for the receiver noise in channels 1 and 2 are  $R_1$  and  $R'_1$  and  $R_2$  and  $R'_2$ , respectively. These amplitudes are related to the temperatures by

$$T_{\mathcal{A}} = A \times A^* + (A') \times (A')^* \tag{2a}$$

$$T_B = B \times B^* + (B') \times (B')^*$$
(2b)

$$T_{R1} = R_1 \times R_1^* + (R_1) \times (R_1)^*$$
(2c)

$$T_{R2} = R_2 \times R_2^* + (R_2) \times (R_2)^*.$$
 (2d)

In general, the 90° beam-splitter hybrid does not split the signals equally but has amplitudes of  $\rho$  for reflection and  $\tau$  for transmission, which are assumed to be the same for both sidebands. Depending on how the LO is injected, the continuous comparison radiometer can be realized with a 180 or 90° input hybrid. After going through the quasioptical hybrid, the signal amplitudes are

$$V_{1} = j \exp(j\omega_{LO}t) \{ 1 + (\rho A - j\tau B + R_{1}) \exp(j\omega_{1}t) + (\rho A' - j\tau B' + R'_{1}) \exp(-j\omega_{1}t) \}$$
(3a)  
$$V_{2} = \exp(j\omega_{LO}t) \{ 1 + (\tau A + j\rho B + R_{2}) \exp(j\omega_{1}t) + (\tau A' + j\rho B' + R'_{2}) \exp(-j\omega_{1}t) \}$$
(3b)

where j is the square root of (-1),  $\omega_{LO}$  the angular frequency of the 97-GHz LO,  $\omega_1$  the first IF angular frequency, and t is time. The signal level of the 97-GHz LO is represented by the first term in the expressions for  $V_1$  and  $V_2$ . It is arbitrarily set to 1 to show the phase of the LO. The second and third terms are the upper and lower sidebands, respectively.

After double-sideband mixing with the 97-GHz LO, gain at the first IF and LSB mixing from 4.7 ( $\omega_1$ ) to 2.1 GHz ( $\omega_2$ ) with a 6.8-GHz second LO, the voltages are

$$V_{3} = \exp(j\phi)G_{1}\{(\rho A - j\tau B + R_{1})\exp(-j\omega_{2}t) + (\rho A' - j\tau B' + R'_{1})\exp(j\omega_{2}t)\} \quad (4a)$$

$$V_{4} = G_{2}\{(\tau A + j\rho B + R_{2})\exp(-j\omega_{2}t) + (\tau A' + j\rho B' + R'_{2})\exp(j\omega_{2}t)\} \quad (4b)$$

where the term  $\exp(j\phi)$  is the relative phase shift introduced by the 6.8-GHz phase shifter and  $G_1$  and  $G_2$  are the voltage gains in the two channels.

At the output of the precision 180° IF hybrid, the voltages are

$$V_5 = (V_3 + V_4)\sqrt{2}$$
 (5a)

$$V_6 = (V_3 - V_4)\sqrt{2}.$$
 (5b)

The detected power is obtained by multiplying  $V_5$  and  $V_6$  by their complex conjugates and low-pass filtering to remove harmonics of the IF components. Cross products such as  $A \times B$  and  $R_1 \times R_2$  will average out to zero when the inputs are uncorrelated and there is no correlated noise in the two receiver temperatures. In the following discussion, K and  $K(1-\delta)$  represent the product of the detector sensitivities and difference amplifier gains for the two detected outputs, where the magnitude of  $(\delta)$  is less than 0.05. In this case, the voltage, of the two outputs are

$$V_{7} = [K/2] \left\{ T_{A} \left[ (\rho G_{1})^{2} + (\tau G_{2})^{2} + (2\rho \tau G_{1}G_{2})\cos(\phi) \right] + T_{B} \left[ (\tau G_{1})^{2} + (\rho G_{2})^{2} - (2\rho \tau G_{1}G_{2})\cos(\phi) \right] + T_{R1} (G_{1})^{2} + T_{R2} (G_{2})^{2}$$
(6a)  
$$V_{8} = [K(1-\delta)/2]$$

$$\cdot \left\{ T_{\mathcal{A}} \Big[ (\rho G_1)^2 + (\tau G_2)^2 - (2\rho \tau G_1 G_2) \cos(\phi) \Big] \\ + T_{\mathcal{B}} \Big[ (\tau G_1)^2 + (\rho G_2)^2 + (2\rho \tau G_1 G_2) \cos(\phi) \Big] \\ + T_{\mathcal{R}1} (G_1)^2 + T_{\mathcal{R}2} (G_2)^2 \right\}.$$
 (6b)

Then, the output of the radiometer is the difference of  $V_7$  and  $V_8$ 

$$V_{\text{out}} = K \{ [(2\rho\tau G_1 G_2)(T_A - T_B)\cos(\phi)] + (\delta/2) \{ T_A [(\rho G_1)^2 + (\tau G_2)^2] + T_B [(\tau G_1)^2 + (\rho G_2)^2)] + T_{R1} (G_1)^2 + T_{R2} (G_2)^2 \}.$$
(7)

The quartz beam splitter splits the signal with 54 percent of the incident power being reflected and 46 percent transmitted. This gives values for  $\rho$  and  $\tau$  of 0.750 and 0.667, so that  $(2\rho\tau)$  equals 0.996. When the two IF gains are equal, as is the case for normal operation, the system output is proportional to the input temperature difference plus a fractional offset term due to the detector imbalance. In this case, the output is given by

$$V_{\text{out}} = (KG_1G_2)\{(T_A - T_B)\cos(\phi) + (\delta)[(T_A + T_B)/2 + (T_{R1} + T_{R2})/2]\}.$$
 (8)

h the precision IF hybrid and careful matching of the detectors,  $\delta$  can be reduced to less than 0.001. This will give an offset on the order of 0.5 K, which can be measured by cycling through the second LO phase  $\phi$ , or can be subtracted out by observing the source alternately in beam A and beam B.

#### III. FRONT-END SYSTEM

The two input beams are 51 mm apart at the Cassegrain focus of a 13.7-m-diam telescope with a f/d ratio of 4.0. This gives two 1.0'-diam beams which are 3.3' apart in azimuth on the sky. This close beam separation was chosen to optimize the cancellation of atmospheric effects. A wider beam separation could be designed if broader sources were to be mapped. The inputs are linearly polarized with the two polarizations 110° apart.

As is schematically shown in Fig. 2, the front-end portion of the system consists of the input optics, the scalar feeds, the cryogenic millimeter mixers, and the IF amplifiers. All of these are integrated into the cryogenics Dewar which is  $250 \times 300 \times 400$  mm in size. The top and bottom covers have O-ring and RF shielding grooves. This design has given a very reliable system. The cooling to 15 and 77 K is done with a CTI 350CP closed-cycle helium refrigerator. A 0.8-mm-thick aluminum shield is attached to the 77

k station of the refrigerator to minimize the thermal radiation loading on the 15 K components. Part of this heat shield is lined with a microwave absorber to act as a black body at 77 K for the LO signal which is not coupled into the mixers and to terminate any reflections and spillover in the optics at 77 K.

As is described in detail in the next section, the optics take the beams at the Cassegrain focus of the antenna, expand them so that the LO can be injected, and combine the two beams in the quasi-optical 90° hybrid. The beams are then refocused to match into the scalar feeds.

The 97.3-GHz LO is provided by a Gunn-diode oscillator having 10 mW of power. The output of the oscillator is isolated and attenuated before going through a waveguide vacuum feedthrough. Inside the Dewar, the LO is coupled into the optics with a rectangular horn, a rexolite focusing lens, and two flat mirrors.

Double-sideband mixing is done in a pair of Schottkydiode millimeter mixers which are cooled to 15 K. The mixers have been developed at the FCRAO and have broad-band RF filters and noncontacting backshorts. They utilize GaAs diodes which have a low doping of  $3 \times 10^{16}$ cm<sup>-3</sup> to give an optimum performance when cryogenically cooled [8], [9]. The mixer blocks are machined from OFHC pper to minimize their RF losses and have a linear taper

om full- to 1/4-height waveguide. The RF filter has been - signed to present a short to the diode at 97 GHz and is reactive at the second harmonic of the LO to minimize



Fig. 2. Radiometer optics diagram. The symmetrical system uses fusedsilica lenses to focus the inputs from the Cassegrain focus into the cryogenic Dewar. A thin sheet of fused-silica acts as a 90° hybrid. Then the 97-GHz LO is injected via a dielectric Fabry-Perot and the beams are refocused into the millimeter mixers.

conversion to higher harmonics. The noncontacting backshort has also been especially designed [10] to be a short circuit at 97 GHz and a pure reactance at 194 GHz. The backshort position was optimized at room temperature and locked into place with a setscrew before cooling the system to 15 K. The whisker length has been optimized since its length tunes the mixer response versus the input frequency [8].

Cooled circulators and FET amplifiers give an IF noise temperature of 25–30 K over the 4.4–5.0-GHz band. Each IF chain consists of an input isolator [12], FET amplifier [13], and output isolator at 20 K, providing a net gain of 11 dB. A second FET with 12 dB more gain is mounted on the 77 K station of the closed-cycle helium refrigerator. The circulators and amplifiers which were made at the FCRAO are followed by commercial amplifiers at 300 K, which are mounted within the vacuum chamber. Their 55 dB of gain gives a net RF to IF gain of 80 dB before the signals leave the RF-shielded Dewar. These signals are then processed as discussed in Section V.

#### IV. OPTICS

The input optics to the radiometer serve multiple purposes. They

- 1) illuminate the subreflector of the 14-m-diam telescope with a 12-dB edge taper,
- 2) inject the local oscillator power at 97.3 GHz,
- combine the input beams in such a way as to create a quasi-optical 90° hybrid for the two sidebands at 93 and 102 GHz.

The optics block diagram is shown in Fig. 2. The photograph in Fig. 3 shows the front end with the vacuum Dewar removed. This view shows the optics after the fused-silica lenses. Fig. 4 shows the beam propagation in one plane so that the various optical elements and the Gaussian beam can be accurately displayed. In the actual

$$IF$$

$$H-12CH_{1}$$

$$V_{1}$$

$$I(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$H(r)^{*}$$

$$F_{1}(h_{n}, h_{n}, h_{n})$$

$$H(r)^{*}$$

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Figure 5.3.4(a). Block diagram of a balanced SIS mixer.



and a second transmission of the second se noise in me unwanted sideband. The schematic is shown chip will be about 2 X 2 mm in size for 200-300 GHz.

gure 5.3.5. We expect that the mixer



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